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LETTER TO THE EDITOR

## Discrete gradient methods for solving ODEs numerically while preserving a first integral

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**Abstract.** We show that all direct methods for preserving a first integral during the numerical integration of an ordinary differential equation fit into the unified framework of discrete gradient methods. Using this framework we construct several new integral-preserving schemes.

### 1. Introduction

In the numerical analysis of dynamical systems there has been an increased emphasis in recent years on qualitative aspects [1, 2]. This has come about for two reasons: first, the ever increasing power of computers that is making feasible the implementation of sophisticated algorithms, hitherto computationally too expensive; second, the discovery that certain *qualitative* features of ODEs can be preserved **exactly** in numerical integration schemes (while simultaneously quantitative features are not neglected)‡. Examples of such qualitative features that can be preserved include Hamiltonian structure [4], the presence of symmetries [5, 6], conservation of phase-space volume in source-free systems [7, 8], the presence of attractors and other  $\omega$ -limit sets [1], and conservation of first integrals [9–19].

In this letter we shall concentrate on the preservation of a single first integral§. More precisely, we shall study a system of coupled autonomous first-order ODEs

$$\frac{dx}{dt} = f(x) \tag{1}$$

where the vector field  $f : \mathbf{R}^n \rightarrow \mathbf{R}^n$  is such that there exists a scalar integral  $I$ , i.e.

$$\frac{dI(x)}{dt} = 0. \tag{2}$$

We wish to find a discrete approximation to (1):

$$\frac{x' - x}{\tau} = \varphi(x, x', \tau) \tag{3}$$

such that the integral  $I$  is preserved exactly, i.e.

$$I(x') = I(x). \tag{4}$$

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‡ These two points are exemplified by the recent 100 million year integration of the equations of motion of the entire solar system, using a symplectic integrator [3].

§ The case of more than one first integral will be treated in [17].

(In (3) and (4) and below, by a slight abuse of notation,  $x = x(n\tau)$  and  $x' = x((n+1)\tau)$ , where  $\tau$  is the time step).

Preserving first integrals is important because of their physical relevance, e.g. in mechanics and astronomy, and also because they can ensure long-term stabilizing effects. They also play a significant role in determining which bifurcations occur generically. It is particularly important that a first integral be preserved if the dimension of the system is low, if the integration time is very long, or if the integral surface is compact.

To date, direct methods for preserving a first integral have mostly fallen into three categories<sup>†</sup>:

- (i) *Integrators which preserve quadratic integrals.* Cooper [9] showed that the implicit midpoint rule

$$\frac{x' - x}{\tau} = f\left(\frac{x' + x}{2}\right) \quad (5)$$

preserves quadratic integrals (actually, so do all symplectic Runge–Kutta methods).

- (ii) *Hamiltonian-preserving methods.* Consider a system of Hamiltonian ODEs,

$$\frac{dx}{dt} = \omega \cdot \frac{DH(x)}{Dx} \quad (6)$$

where  $\omega$  is the standard symplectic structure

$$\omega := \begin{pmatrix} 0 & \text{Id} \\ -\text{Id} & 0 \end{pmatrix} \quad (7)$$

(with Id the  $n/2$ -dimensional unit matrix), and the gradient is denoted by

$$\frac{DH(x)}{Dx} := \begin{pmatrix} \partial H(x)/\partial x_1 \\ \vdots \\ \partial H(x)/\partial x_n \end{pmatrix}. \quad (8)$$

Some Hamiltonian-preserving methods for (6) are given in [10–14, 19].

- (iii) *Splitting methods which preserve general integrals.* Quispel and Capel [16] found a direct method for preserving general integrals. It consisted of four steps:

- Write the ODE in skew-gradient form

$$\frac{dx}{dt} = S(x) \cdot \frac{DI(x)}{Dx} \quad (9)$$

where  $S$  is some skew-symmetric matrix (i.e.  $S^t = -S$ ).

- Split the right-hand side of (9) into two-dimensional vector fields, each possessing the integral  $I$ .
- Integrate these two-dimensional vector fields using an essentially two-dimensional integral-preserving integrator.
- Construct an  $n$ -dimensional integral-preserving integrator from the two-dimensional integrators obtained in the previous step.

<sup>†</sup> Some indirect methods are also given in [15, 20].

## 2. Necessary and sufficient conditions for a dynamical system to possess a first integral

In [16] we proved the following theorem:

*Theorem 1.* The autonomous ODE

$$\frac{dx}{dt} = f(x) \quad (10)$$

has a first integral  $I(x)$  if and only if (10) can (formally) be written as a skew-gradient system:

$$\frac{dx}{dt} = S(x) \cdot \frac{DI}{Dx}(x). \quad (11)$$

□

In (11),  $S$  is some skew-symmetric matrix

$$S^t = -S \quad (12)$$

and the gradient is denoted by

$$\frac{DI}{Dx}(x) := \begin{pmatrix} \partial I(x)/\partial x_1 \\ \vdots \\ \partial I(x)/\partial x_n \end{pmatrix}. \quad (13)$$

The (constructive) proof of theorem 1 is given in [16]. It uses the infinitesimal identity

$$dx \cdot \frac{DI}{Dx} = dI. \quad (14)$$

It turns out that theorem 1 carries over to the discrete case:

*Theorem 2.* Let the ‘discrete gradient’  $\Delta I/\Delta x$  be any vector satisfying the crucial discrete analogue of identity (14) [10, 12, 19]:

$$(x' - x) \cdot \frac{\Delta I}{\Delta x}(x, x') = I(x') - I(x). \quad (15)$$

Then the (implicit) mapping

$$\frac{x' - x}{\tau} = \varphi(x, x', \tau) \quad (16)$$

has a first integral  $I(x)$  if and only if (16) can (formally) be written as a discrete skew-gradient system

$$\frac{x' - x}{\tau} = \tilde{S}(x, x', \tau) \cdot \frac{\Delta I}{\Delta x}(x, x'). \quad (17)$$

□

In (17),  $\tilde{S}$  is some skew-symmetric matrix

$$\tilde{S}^t = -\tilde{S}. \quad (18)$$

The proof is analogous to the proof of theorem 1, and will therefore be omitted. Note, however, that for a given integral the choice of discrete gradient satisfying (15) is not unique, and that for a given discrete gradient the matrix  $\tilde{S}$  is not unique (see below).

### 3. Discrete gradient methods as integral-preserving integrators

Theorem 2 implies that the exact flow of an integral-possessing ODE is a discrete skew-gradient system. Combining theorems 1 and 2 then suggests the following two-step method for constructing integral-preserving integrators (IPIs):

(i) Write the ODE in the skew-gradient form

$$\frac{dx}{dt} = S(x) \cdot \frac{DI(x)}{Dx}. \quad (19)$$

(ii) Integrate (19), using an integral-preserving discrete gradient method

$$\frac{x' - x}{\tau} = \tilde{S}(x, x', \tau) \cdot \frac{\Delta I}{\Delta x}(x, x'). \quad (20)$$

Step (i): Write the ODE in the form  $dx/dt = S \cdot (DI/Dx)$ .

For a given vector field  $f$  and integral  $I$ , there are in general infinitely many choices for the skew matrix  $S$  (see [16]). Here we restrict ourselves to giving a family of particular solutions. For any arbitrary vector field  $a(x)$ , a skew matrix  $S$  as in theorem 1 is given by [21]

$$S_{i,j}(x) = \frac{f_i(x)a_j(x) - a_i(x)f_j(x)}{\sum_{k=1}^n (a_k(x) \left( \partial I(x) / \partial x_k \right))}. \quad (21)$$

(This is easily shown using  $f \cdot (DI/Dx) = 0$ .) The choice of the vector field  $a(x)$  is governed by our two conflicting desires for  $S$  to be as simple as possible, but preferably free from singularities.

*Example 1.*  $\exists k, \forall x, (\partial I(x) / \partial x_k) \neq 0$ . In this case we can choose

$$a(x) = \hat{e}_k \quad (22)$$

where  $\hat{e}_k$  is a unit vector in the  $x_k$  direction.

*Example 2. The general case.* In this case we can choose<sup>†</sup>

$$a(x) = \frac{DI(x)}{Dx}. \quad (23)$$

Step (ii): Integral-preserving discrete gradient methods.

Even though, formally, all discrete gradients are equivalent (cf theorem 2), in practice (depending on the vector field) one may be more convenient than another:

(iia) Choice of discrete gradient.

Each discrete gradient must satisfy the discrete identity (6), and must be consistent:

$$\left( \frac{\Delta I}{\Delta x} \right)(x, x') = \frac{DI}{Dx}(x) + O(\tau). \quad (24)$$

<sup>†</sup> A sufficient condition for this to work is that  $DI/Dx \neq 0$  on the integral surface on which we are integrating.

A discrete gradient satisfying (15) and (24) was given by Itoh and Abe [10], cf also [12]:

$$\left(\frac{\Delta I}{\Delta x}\right)_1(x, x') := \begin{pmatrix} \frac{I(x'_1, x_2, x_3, x_4, \dots, x_n) - I(x_1, x_2, x_3, x_4, \dots, x_n)}{x'_1 - x_1} \\ \frac{I(x'_1, x'_2, x_3, x_4, \dots, x_n) - I(x'_1, x_2, x_3, x_4, \dots, x_n)}{x'_2 - x_2} \\ \frac{I(x'_1, x'_2, x'_3, x_4, \dots, x_n) - I(x'_1, x'_2, x_3, x_4, \dots, x_n)}{x'_3 - x_3} \\ \vdots \\ \frac{I(x'_1, x'_2, \dots, x'_{n-1}, x'_n) - I(x'_1, x'_2, \dots, x'_{n-1}, x_n)}{x'_n - x_n} \end{pmatrix}. \quad (25)$$

An equally good discrete gradient is obtained by performing an arbitrary permutation  $\pi$  of the indices  $1, \dots, n$ :

$$\left(\frac{\Delta I}{\Delta(\pi x)}\right)_{(j)}(x, x') := \frac{(E_j - 1)}{x'_j - x_j} \left( \prod_{\ell=1}^{\pi^{-1}(j)-1} E_{\pi(\ell)} \right) I(x_1, \dots, x_n) \quad (26)$$

where the  $\ell$ th time shift operator  $E_\ell$  is defined by

$$E_\ell g(\dots, x_\ell, \dots) := g(\dots, x'_\ell, \dots) \quad (\ell = 1, \dots, n) \quad (27)$$

where  $g$  is any function and the dots in (27) denote the other  $n - 1$  (primed or unprimed) variables. This allows us to construct a quite general discrete gradient:

$$\left(\frac{\Delta I}{\Delta x}\right)_2(x, x') := \frac{\sum_{\pi} c(\pi) (\Delta I / \Delta(\pi x))}{\sum_{\pi} c(\pi)} \quad (28)$$

where the  $c(\pi)$  are arbitrary coefficients. Choosing

$$c(\pi) = \begin{cases} 1 & \text{if } \pi = \text{identity} \\ 0 & \text{otherwise} \end{cases} \quad (29)$$

we recover  $(\Delta I / \Delta x)_1$ . Choosing

$$c(\pi) = \begin{cases} \frac{1}{2} & \text{if } \pi = \text{identity} \\ \frac{1}{2} & \text{if } \pi(1, \dots, n) = (n, \dots, 1) \\ 0 & \text{otherwise} \end{cases} \quad (30)$$

we obtain

$$\left(\frac{\Delta I}{\Delta x}\right)_3(x, x') = \frac{(\Delta I / \Delta x)_1(x, x') + (\Delta I / \Delta x)_1(x', x)}{2}. \quad (31)$$

A more symmetric IPI is obtained by choosing

$$c(\pi) = \frac{1}{n!} \quad \forall \pi \quad (32)$$

leading to

$$\left(\frac{\Delta I}{\Delta x}\right)_4(x, x') = \frac{1}{n!} \sum_{\pi} \frac{\Delta I}{\Delta(\pi x)}(x, x'). \quad (33)$$

Note that any discrete gradient can be obtained from any other by adding a solution of the (homogeneous) equations

$$(x' - x)a(x, x') = 0 \quad a(x, x') = \mathcal{O}(\tau). \quad (34)$$

The solution to (34) is

$$a(x, x') = b(x, x') - (x' - x) \frac{(x' - x) \cdot b(x, x')}{(x' - x) \cdot (x' - x)} \quad (35)$$

where  $b$  is an arbitrary vector of order  $\tau$ .

We choose

$$b(x, x') = \frac{DI}{Dx}(z(x, x')) - \left(\frac{\Delta I}{\Delta x}\right)_1(x, x') \quad (36)$$

where  $z(x, x') = x + \mathcal{O}(\tau)$ . Substituting in (35) and adding to (25) we obtain

$$\left(\frac{\Delta I}{\Delta x}\right)_5(x, x') = \frac{DI}{Dx}(z) - (x' - x) \frac{(x' - x) \cdot (DI/Dx)(z) - I(x') + I(x)}{(x' - x) \cdot (x' - x)}. \quad (37)$$

This very nice solution was discovered by Gonzalez.

(iib) Choice of skew-symmetric matrix.

To be consistent each skew matrix must satisfy

$$\tilde{S}(x, x', \tau) = S(x) + \mathcal{O}(\tau). \quad (38)$$

Two simple choices are therefore

$$\tilde{S}_1(x, x', \tau) = S(x) \quad (39)$$

and

$$\tilde{S}_2(x, x', \tau) = S\left(\frac{x + x'}{2}\right). \quad (40)$$

(iic) Combination of the discrete gradient and the skew matrix.

A simple first-order IPI is

$$\frac{x' - x}{\tau} = S(x) \left(\frac{\Delta I}{\Delta x}\right)_1(x, x'). \quad (41)$$

To obtain a second-order IPI we impose the additional requirement of 'time symmetry', i.e.

$$\tilde{S}(x, x', \tau) = \tilde{S}(x', x, -\tau) \quad (42)$$

$$\left(\frac{\Delta I}{\Delta x}\right)(x, x') = \left(\frac{\Delta I}{\Delta x}\right)(x', x). \quad (43)$$

Equation (42) is satisfied by  $\tilde{S}_2$ , and (43) is satisfied e.g. by  $(\Delta I/\Delta x)_3$ ,  $(\Delta I/\Delta x)_4$ , and, if  $z(x, x') = z(x', x)$ , by  $(\Delta I/\Delta x)_5$  [19].

The 'time-symmetry' of (42) and (43) makes these integrators suitable building blocks for constructing integrators of arbitrary order using Yoshida's method [22].

**4. A numerical example**

To illustrate the fact that standard methods [23] do not, in general, preserve first integrals, we have compared the performance of three numerical integration methods on the ODE

$$\begin{aligned} \dot{x} &= x + yz + 0.1xy^2 + xy^3 + 0.1xy^5 \\ \dot{y} &= -x^2 + z^2 - 0.1x^2y^2 \\ \dot{z} &= -z - xy - y^3z \end{aligned} \tag{44}$$

with initial conditions  $x(0) = z(0) = 1, y(0) = 2$ . Note that the system (44) has the first integral

$$I = \frac{x^2}{2} + \frac{y^4}{4} + y + \frac{z^2}{2}. \tag{45}$$

The level surfaces of this integral are all compact. The three fourth-order methods we used are an explicit linear multi-step method (LM4), an explicit Runge–Kutta method (RK4), and one of our new implicit IPIs (QT4). The time steps  $\tau$  have been adjusted such that the amount of numerical work performed by the three methods is identical.

- LM4: This is the fourth-order Adams–Bashforth method.
- RK4: The standard fourth-order Runge–Kutta method defined by the Butcher tableau:

$$\begin{array}{c|ccc} \frac{1}{2} & & & \\ 0 & \frac{1}{2} & & \\ 0 & 0 & 1 & \\ \frac{1}{6} & \frac{2}{6} & \frac{2}{6} & \frac{1}{6} \end{array} \tag{46}$$

- QT4: Note that (44) can be written in skew-gradient form (11), with

$$S = \begin{pmatrix} 0 & x + 0.1xy^2 & y \\ -x - 0.1xy^2 & 0 & z \\ -y & -z & 0 \end{pmatrix}. \tag{47}$$

QT4 was obtained by first constructing a second-order IPI, using (17) with (31) and (43), and then applying Yoshida’s method to obtain a fourth-order IPI [22].

Numerical results are given in figure 1.

**5. Reduction of the discrete gradient method to previous methods**

- (i) *Quadratic integrals: reduction to the implicit midpoint rule.* From the arguments given above it follows that there formally exists a matrix  $S$  such that

$$f(x) = S(x) \cdot \frac{DI(x)}{Dx}. \tag{48}$$

Now use an IPI with  $\tilde{S}_2$  and e.g.  $(\Delta I/\Delta x)_4$  or  $(\Delta I/\Delta x)_5$  (with  $z = (x + x')/2$ ). It can be shown that if  $I$  is quadratic, then

$$\left(\frac{\Delta I}{\Delta x}\right)_4(x, x') = \left(\frac{\Delta I}{\Delta x}\right)_5(x, x') = \frac{DI}{Dx}\left(\frac{x + x'}{2}\right) \tag{49}$$

and using (48) we see this reduces to the implicit midpoint rule (5). By the same token, any IPI with (40) preserves all quadratic Casimirs.



- (ii) *Hamiltonian systems: reduction to the Itoh–Abe and Gonzalez methods.* For standard Hamiltonian systems we can take  $I = H$  and  $S = \tilde{S} = \omega$ . The Itoh–Abe method is then recovered choosing the discrete gradient  $(\Delta H/\Delta x)_1$ . Gonzalez’s method is obtained choosing the discrete gradient  $(\Delta I/\Delta x)_5$  with  $z = (x + x')/2$ .
- (iii) *Reduction to integral-preserving splitting methods.* The ODE (11) is equivalent to the vector field

$$v = \sum_{i < j} v_{i,j} \quad (50)$$

where

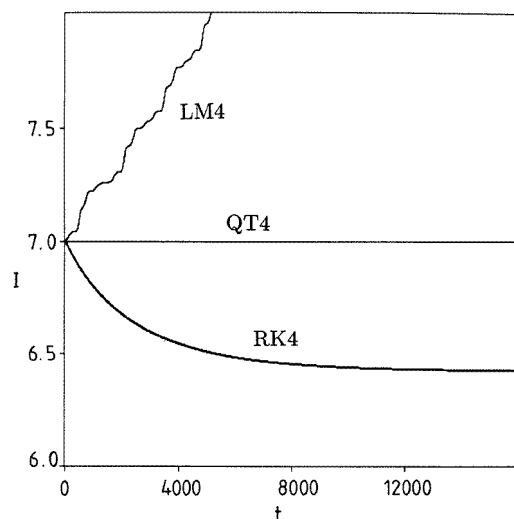
$$v_{i,j} := S_{i,j} \left( \frac{\partial I}{\partial x_j} \frac{\partial}{\partial x_i} - \frac{\partial I}{\partial x_i} \frac{\partial}{\partial x_j} \right) \quad (51)$$

and each vector field  $v_{i,j}$  preserves the integral  $I$ . Using the splitting (50) we obtain the method of Quispel and Capel [16]. Using various other partitions of the set  $\{v_{i,j}\}$ , there are many ways in which the  $n$ -dimensional vector field  $v$  can be split into lower-dimensional vector fields that each preserve  $I$ .

## 6. Concluding remark

Associated with the continuous system (11) there is a bracket formulation

$$\frac{dh(x)}{dt} = \{h, I\}_C \quad (52)$$



**Figure 1.** This figure illustrates that conventional integration methods (RK4 and LM4) do not preserve first integrals. It shows the numerical value of the integral  $I$ , as a function of time  $t = n\tau$ , for the three-dimensional test system (44). It compares the standard fourth-order Runge–Kutta method (RK4;  $\tau = 0.031$ ), a fourth-order linear multistep method (LM4;  $\tau = 0.040$ ) and a fourth-order discrete gradient method (QT4;  $\tau = 0.443$ ). (The time steps  $\tau$  have been adjusted so that the amount of numerical work performed by the three methods is equivalent.)

where the continuous bracket is defined by

$$\{h, g\}_C := \frac{Dh}{Dx} \cdot S \cdot \frac{Dg}{Dx}. \quad (53)$$

If  $\{h, I\}_C \equiv 0$ , then  $h$  is a first integral.

Associated with the discrete system (17) there is also a bracket formulation

$$\frac{h(x') - h(x)}{\tau} = \{h, I\}_D \quad (54)$$

where the discrete bracket is given by

$$\{h, g\}_D := \frac{\Delta h}{\Delta x} \cdot \tilde{S} \cdot \frac{\Delta g}{\Delta x}. \quad (55)$$

If  $\{h, I\}_D \equiv 0$ , then  $h$  is a first integral.

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## References

- [1] Stuart A M 1994 Numerical analysis of dynamical systems *Acta Numerica* ed A Iserles
- [2] Iserles A and Zanna A 1995 *Qualitative numerical analysis of ODEs* DAMTP Tech. Rep., Cambridge University
- [3] Sussman G J and Wisdom J 1992 *Science* **257** 56
- [4] Sanz-Serna M and Calvo M P 1994 *Numerical Hamiltonian Problems* (London: Chapman and Hall)
- [5] Stoffer D 1995 *Computing* **55** 1
- [6] McLachlan R I, Quispel G R W and Turner G S 1996 *Numerical integrators that preserve symmetries and reversing symmetries* Latrobe University Preprint
- [7] Feng Kang and Wang Dao-Liu 1994 Dynamical systems and geometric construction of algorithms *Contemporary Mathematics* **163** ed Z Shi and C Yang
- [8] Quispel G R W 1995 *Phys. Lett.* **206A** 26
- [9] Cooper G J 1987 *IMA J. Num. Anal.* **7** 1
- [10] Itoh T and Abe K 1988 *J. Comput. Phys.* **77** 85
- [11] Greenspan D 1995 *Comp. Math. Appl.* **27** 37
- [12] Kriksin Yu A 1993 *Zh. Vychisl. Mat. Mat. Fiz.* **33** 206
- [13] Chorin A J, Hughes T J R, McCracken M F and Marsden J E 1978 *Comm. Pure Appl. Maths.* **XXXI** 205
- [14] Marsden J E 1992 *Lectures on Mechanics* (Cambridge: Cambridge University Press)
- [15] Gear C W 1992 *Physica D* **60** 303
- [16] Quispel G R W and Capel H W 1996 *Solving ODEs numerically while preserving a first integral* to appear in *Phys. Lett. A*
- [17] Quispel G R W, Capel H W and Scovel C 1996 *Discrete gradient methods for solving ODEs numerically while preserving all first integrals* in preparation
- [18] Calvo M P, Iserles A and Zanna A 1995 *Runge–Kutta methods on manifolds* DAMTP Tech. Rep., Cambridge University
- [19] Gonzalez O 1996 *Time Integration and Discrete Hamiltonian Systems* Preprint
- [20] Shampine L F 1986 *Comp. Math. Appl.* B **12** 1278
- [21] A special instance of this is given in: Hojman S A 1996 *J. Phys. A* **29** 667
- [22] Yoshida H 1990 *Phys. Lett.* **150A** 262
- [23] Hairer E, Nørsett S P and Wanner G 1992 *Solving Ordinary Differential Equations* (Berlin: Springer)